

NAPREDNA KVANTNA MEHANIKA

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ZADATAK 1

Operatore stvaranja i poništenja za *bozone* definiramo na sljedeći način:

$$a_i^\dagger |n_1 n_2 \dots n_i \dots n_\infty\rangle = \sqrt{n_i + 1} |n_1 n_2 \dots n_i + 1, \dots n_\infty\rangle$$

$$a_i |n_1 n_2 \dots n_i \dots n_\infty\rangle = \sqrt{n_i} |n_1 n_2 \dots n_i - 1, \dots n_\infty\rangle$$

Pokažite da za operatore stvaranja i poništenja vrijede komutacijske relacije

$$[a_k, a_l^\dagger] = a_k a_l^\dagger - a_l^\dagger a_k = \delta_{kl}$$

$$[a_k, a_l] = [a_k^\dagger, a_l^\dagger] = 0$$

ZADATAK 2

(a) Na vježbama smo spomenuli da operator

$$\boldsymbol{\Sigma} \cdot \hat{\mathbf{n}}$$

gdje je $\hat{\mathbf{n}}$ proizvoljan jedinični vektor, ne komutira s Diracovim hamiltonijanom. Nasuprot tome, operator heliciteta

$$\boldsymbol{\Sigma} \cdot \hat{\mathbf{p}}$$

komutira s Diracovim hamiltonijanom, gdje je $\hat{\mathbf{p}}$ jedinični vektor u smjeru operatora impulsa. Dokažite ovu tvrdnju.

(b) Pokažite da su ± 1 svojstvene vrijednosti operatora heliciteta.

(c) Pokažite da je operator heliciteta rotacijski invarijantan.

1. Neka je $l > k$. Vektori baze u Fockovom prostoru su $\{|n_1 n_2 \dots n_\infty\rangle\}$.

$$a_k a_e^\dagger |n_1 n_2 \dots n_\infty\rangle = a_k (\sqrt{n_{e+1}} |n_1, \dots, n_{e+1}, \dots, n_\infty\rangle) \\ = \sqrt{n_{e+1}} \sqrt{n_k} |n_1, \dots, n_{k-1}, \dots, n_{e+1}, \dots, n_\infty\rangle$$

$$a_e^\dagger a_k |n_1 n_2 \dots n_\infty\rangle = a_e^\dagger (\sqrt{n_k} |n_1, \dots, n_{k-1}, \dots, n_\infty\rangle) \\ = \sqrt{n_k} \sqrt{n_{e+1}} |n_1, \dots, n_{k-1}, \dots, n_{e+1}, \dots, n_\infty\rangle$$

Budući je $l > k$, $l \neq k$ imamo

$$(a_k a_e^\dagger - a_e^\dagger a_k) |n_1 n_2 \dots n_\infty\rangle = 0$$

Neka je $l = k$. Tada je

$$a_k a_k^\dagger |n_1 \dots n_\infty\rangle = a_k (\sqrt{n_{k+1}} |n_1, \dots, n_{k+1}, \dots, n_\infty\rangle) \\ = \underbrace{\sqrt{n_{k+1}} \sqrt{n_{k+1}}}_{(n_{k+1})} |n_1, \dots, n_k, \dots, n_\infty\rangle$$

$$a_k^\dagger a_k |n_1 \dots n_\infty\rangle = a_k^\dagger (\sqrt{n_k} |n_1, \dots, n_{k-1}, \dots, n_\infty\rangle) \\ = \underbrace{\sqrt{n_k} \sqrt{n_k}}_{n_k} |n_1, \dots, n_k, \dots, n_\infty\rangle$$

$$(a_k a_k^\dagger - a_k^\dagger a_k) |n_1 n_2 \dots n_\infty\rangle = |n_1 n_2 \dots n_\infty\rangle$$

Prema tome,

$$[a_k, a_e^\dagger] |n_1 n_2 \dots n_\infty\rangle = \delta_{ke} |n_1 n_2 \dots n_\infty\rangle$$

Proizvoljni vektor stanja je linearna kombinacija stanja

$\{|n_1 n_2 \dots n_\infty\rangle\}$. Prema tome, vrijedi

$$\begin{aligned} [a_k, a_e^\dagger] |\psi\rangle &= [a_k, a_e^\dagger] \sum_{\{n_i\}} c_{\{n_i\}} |n_1 n_2 \dots n_\infty\rangle \\ &= \sum_{\{n_i\}} c_{\{n_i\}} [a_k, a_e^\dagger] |n_1 n_2 \dots n_\infty\rangle \\ &= \delta_{ke} |\psi\rangle \end{aligned}$$

Budući je $|\psi\rangle$ proizvoljan vektor stanja

$$[a_k, a_e^\dagger] = \delta_{ke}$$

Provjerimo koliko je $[a_k, a_e]$

$$a_k a_e |n_1 n_2 \dots n_\infty\rangle = a_k (\sqrt{n_e} |n_1 \dots, n_e - 1, \dots, n_\infty\rangle)$$

$$= \sqrt{n_e} \sqrt{n_k} |n_1 \dots, n_k - 1, \dots, n_e - 1, \dots, n_\infty\rangle$$

$$a_e a_k |n_1 n_2 \dots n_\infty\rangle = \sqrt{n_e} \sqrt{n_k} |n_1 \dots, n_k - 1, \dots, n_e - 1, \dots, n_\infty\rangle$$

$$(a_k a_e - a_e a_k) |n_1 n_2 \dots n_\infty\rangle = 0$$

Čak i ako je $e = k$,

$$a_k^2 |n_1 n_2 \dots n_\infty\rangle = \sqrt{n_k} \sqrt{n_k - 1} |n_1 \dots, n_k - 2, \dots, n_\infty\rangle$$

pa je

$$(a_k^2 - a_k^2) |n_1 n_2 \dots n_\infty\rangle = 0$$

Prema tome,

$$[a_k, a_e] = 0$$

Ista logika vrijedi i za $[a_k^+, a_e^+] = 0$.

2.

(a) Diracov hamiltonijan

$$H = c \vec{\alpha} \cdot \vec{p} + mc^2 \beta, \quad \vec{p} \text{ je operator}$$

$$\alpha_j = -\gamma_5 \Sigma_j = -\Sigma_j \gamma_5 \quad (\text{to vještbi})$$

$$\gamma_5 = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix}$$

$$\Sigma_k = \begin{pmatrix} \sigma_k & 0 \\ 0 & \sigma_k \end{pmatrix}$$

$$\vec{\alpha} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix}; \quad \vec{\beta} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$[H, \vec{\Sigma} \cdot \vec{n}] = [c \alpha_j p_j + mc^2 \beta, \sum_k n_k \Sigma_k]$$

j, k, m nijemi indeksi; sumiraju se po njima

$$[\beta, \Sigma_k] = \beta \Sigma_k - \Sigma_k \beta$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} \sigma_k & 0 \\ 0 & \sigma_k \end{pmatrix} - \begin{pmatrix} \sigma_k & 0 \\ 0 & \sigma_k \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$= \begin{pmatrix} \sigma_k & 0 \\ 0 & -\sigma_k \end{pmatrix} - \begin{pmatrix} \sigma_k & 0 \\ 0 & -\sigma_k \end{pmatrix} = 0$$

komutiraju! Ostaje,

$$\begin{aligned} [\alpha_j p_j, \sum_k n_k \Sigma_k] &= [\alpha_j, \sum_k \Sigma_k] p_j n_k \\ &= -\gamma_5 [\sum_j \Sigma_j, \sum_k \Sigma_k] p_j n_k \end{aligned}$$

$$[\Sigma_j, \Sigma_k] = 2i \epsilon_{jke} \Sigma_e$$

Isti uvjet zadovoljavamo \vec{b} uvek! Imamo

$$\begin{aligned} -c \delta_5 \sum_e \underbrace{2i \epsilon_{jke}}_{-de} P_j n_k &= 2i \underbrace{\epsilon_{jke}}_{\epsilon_{kej}} c \alpha_e P_j n_k \\ &= 2ic \vec{n} \cdot (\vec{d} \times \vec{p}) \end{aligned}$$

Samo ako je $\vec{n} \equiv \vec{p}$ ako je uvek. U protivnom, nije!

(b) Operator heliciteta

$$\vec{\Sigma} \cdot \vec{p} = \sum_k \hat{p}_k = \begin{pmatrix} \vec{b} \cdot \vec{p} & 0 \\ 0 & \vec{b} \cdot \vec{p} \end{pmatrix}$$

Tražimo

$$\det \begin{pmatrix} \vec{b} \cdot \vec{p} - \lambda 1 & 0 \\ 0 & \vec{b} \cdot \vec{p} - \lambda 1 \end{pmatrix} = 0$$

Dovoljno je izračunati

$$\det(\vec{b} \cdot \vec{p} - \lambda 1) = 0$$

pa je

$$\det \begin{pmatrix} A & 0 \\ 0 & B \end{pmatrix} = \det A \cdot \det B$$

za blok matricu. Imamo

$$\det \begin{pmatrix} p_3 - \lambda & p_1 - ip_2 \\ p_1 + ip_2 & -p_3 - \lambda \end{pmatrix} = 0$$

$$-(p_3 + \lambda)(p_3 - \lambda) - (p_1 - ip_2)(p_1 + ip_2) = 0$$

$$-(p_3^2 - \Lambda^2) - (p_1^2 + p_2^2) = 0$$

$$\Lambda^2 - (p_1^2 + p_2^2 + p_3^2) = 0$$

No, \vec{p} je poljubni vektor za koreq je

$$p_1^2 + p_2^2 + p_3^2 = 1$$

Prema tome,

$$\Lambda^2 = 1 \Rightarrow \Lambda = \pm 1$$

(c) Operator rotacije je oblika za česticu spina $\frac{1}{2}$ oko osi \hat{n} za kut ϕ

$$\begin{aligned} \mathcal{D}(\hat{n}, \phi) &= \exp\left(-\frac{i \vec{\sigma} \cdot \hat{n}}{2} \phi\right) \\ &= \cos\left(\frac{\phi}{2}\right) - i \vec{\sigma} \cdot \hat{n} \sin\left(\frac{\phi}{2}\right) \end{aligned}$$

Očita generalizacija za 4-dimenzionalni prostor je

$$\begin{aligned} \mathcal{D}(\hat{n}, \phi) &= \exp\left(-\frac{i \vec{\Sigma} \cdot \hat{n}}{2} \phi\right) \\ &= \cos\left(\frac{\phi}{2}\right) - i \vec{\Sigma} \cdot \hat{n} \sin\left(\frac{\phi}{2}\right) \end{aligned}$$

Treba izračunati

$$\begin{aligned} \mathcal{D}^\dagger(\hat{n}, \phi) (\vec{\Sigma} \cdot \hat{p}) \mathcal{D}(\hat{n}, \phi) &= \left(\cos\left(\frac{\phi}{2}\right) + i \vec{\Sigma} \cdot \hat{n} \sin\left(\frac{\phi}{2}\right) \right) \\ &\quad \cdot (\vec{\Sigma} \cdot \hat{p}) \left(\cos\left(\frac{\phi}{2}\right) - i \vec{\Sigma} \cdot \hat{n} \sin\left(\frac{\phi}{2}\right) \right) \\ &= (\vec{\Sigma} \cdot \hat{p}) \cos^2\left(\frac{\phi}{2}\right) - i (\vec{\Sigma} \cdot \hat{p}) (\vec{\Sigma} \cdot \hat{n}) \cos\left(\frac{\phi}{2}\right) \sin\left(\frac{\phi}{2}\right) \\ &\quad + i (\vec{\Sigma} \cdot \hat{n}) (\vec{\Sigma} \cdot \hat{p}) \sin\left(\frac{\phi}{2}\right) \cos\left(\frac{\phi}{2}\right) + (\vec{\Sigma} \cdot \hat{n}) (\vec{\Sigma} \cdot \hat{p}) (\vec{\Sigma} \cdot \hat{n}) \sin^2\left(\frac{\phi}{2}\right) \quad (*) \end{aligned}$$

Koristimo identitet koji vrijedi za $\hat{\sigma}$ -matrice:

$$(\vec{\Sigma} \cdot \vec{p})(\vec{\Sigma} \cdot \vec{n}) = \vec{p} \cdot \vec{n} + i \vec{\Sigma} \cdot (\vec{p} \times \vec{n})$$

$$(\vec{\Sigma} \cdot \vec{n})(\vec{\Sigma} \cdot \vec{p}) = \vec{n} \cdot \vec{p} + i \vec{\Sigma} \cdot (\vec{n} \times \vec{p})$$

$$\vec{p} \cdot \vec{n} = \vec{n} \cdot \vec{p} \quad \text{to su jednaki velični, uvek operativni}$$

Prema stavu, 2. i 3. član u (*) daje

$$\begin{aligned} & \vec{\Sigma} \cdot (\vec{p} \times \vec{n}) \cos\left(\frac{\phi}{2}\right) \sin\left(\frac{\phi}{2}\right) - \vec{\Sigma} \cdot \underbrace{(\vec{n} \times \vec{p})}_{-\vec{p} \times \vec{n}} \cos\left(\frac{\phi}{2}\right) \sin\left(\frac{\phi}{2}\right) \\ &= 2 \vec{\Sigma} \cdot (\vec{p} \times \vec{n}) \cos\left(\frac{\phi}{2}\right) \sin\left(\frac{\phi}{2}\right) \end{aligned}$$

Računamo 4. član

$$\begin{aligned} & (\vec{\Sigma} \cdot \vec{n}) \cdot (\vec{\Sigma} \cdot \vec{p}) (\vec{\Sigma} \cdot \vec{n}) \\ &= (\vec{\Sigma} \cdot \vec{n}) \left(\vec{p} \cdot \vec{n} + i \vec{\Sigma} \cdot (\vec{p} \times \vec{n}) \right) \\ &= (\vec{p} \cdot \vec{n}) \vec{\Sigma} \cdot \vec{n} + i (\vec{\Sigma} \cdot \vec{n}) (\vec{\Sigma} \cdot (\vec{p} \times \vec{n})) \end{aligned}$$

No, po istom pravilu:

$$\begin{aligned} (\vec{\Sigma} \cdot \vec{n}) (\vec{\Sigma} \cdot (\vec{p} \times \vec{n})) &= \underbrace{\vec{n} \cdot (\vec{p} \times \vec{n})}_{=0} + i \vec{\Sigma} \cdot \left[\underbrace{\vec{n} \times (\vec{p} \times \vec{n})}_{\vec{p}(\vec{n} \cdot \vec{n}) - \vec{n}(\vec{n} \cdot \vec{p})} \right] \\ &= i \vec{\Sigma} \cdot \vec{p} - i (\vec{\Sigma} \cdot \vec{n}) (\vec{n} \cdot \vec{p}) \stackrel{=1}{=} \end{aligned}$$

Prema tome,

$$(\vec{\Sigma} \cdot \vec{n})(\vec{\Sigma} \cdot \vec{p})(\vec{\Sigma} \cdot \vec{n}) = 2(\vec{\Sigma} \cdot \vec{n})\vec{p} \cdot \vec{n} - \vec{\Sigma} \cdot \vec{p}$$

Imamo,

$$\begin{aligned} \mathcal{D}^\dagger(\vec{n}, \phi)(\vec{\Sigma} \cdot \vec{p})\mathcal{D}(\vec{n}, \phi) &= (\vec{\Sigma} \cdot \vec{p}) \cos^2 \frac{\phi}{2} \\ &+ 2\vec{\Sigma} \cdot (\vec{p} \times \vec{n}) \cos \frac{\phi}{2} \sin \frac{\phi}{2} \\ &+ 2(\vec{\Sigma} \cdot \vec{n})(\vec{p} \cdot \vec{n}) \sin^2 \frac{\phi}{2} - (\vec{\Sigma} \cdot \vec{p}) \sin^2 \frac{\phi}{2} \end{aligned}$$

$$\cos^2 \frac{\phi}{2} - \sin^2 \frac{\phi}{2} = \cos \phi$$

$$2 \cos \frac{\phi}{2} \sin \frac{\phi}{2} = \sin \phi$$

$$2 \sin^2 \frac{\phi}{2} = 1 - \cos \phi$$

Govorimo izraz postaje

$$\begin{aligned} &= (\vec{\Sigma} \cdot \vec{p}) \cos \phi + \vec{\Sigma} \cdot (\vec{p} \times \vec{n}) \sin \phi + (\vec{\Sigma} \cdot \vec{n})(\vec{p} \cdot \vec{n})(1 - \cos \phi) \\ &= \vec{\Sigma} \cdot \left(\vec{p} \cos \phi + (\vec{p} \times \vec{n}) \sin \phi + \vec{n} (\vec{p} \cdot \vec{n})(1 - \cos \phi) \right) \end{aligned}$$

Iznaz u zagradi je Rodriguesova formula za vrtanje vektora \vec{p} oko \vec{n} za kut ϕ (altimna rotacija)

$$\vec{p}' = \vec{p} \cos \phi + (\vec{p} \times \vec{n}) \sin \phi + \vec{n} (\vec{p} \cdot \vec{n})(1 - \cos \phi)$$

Pokazali smo da vrijedi

$$\mathcal{D}^\dagger(\vec{n}, \phi)(\vec{\Sigma} \cdot \vec{p})\mathcal{D}(\vec{n}, \phi) = \vec{\Sigma} \cdot \vec{p}'$$