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8 Valne funkcije i vektori stanja

8.1 Čestica je u kvantnom stanju opisanom valnom funkcijom

$$\Psi(x) = \begin{cases} \frac{1}{\sqrt{2L}}, & |x| \leq L \\ 0, & |x| > L \end{cases}$$

- (a) Kolika je gustoća vjerojatnosti $\rho(p)$ nalaženja čestice s impulsom p ?
 (b) Nacrtajte gustoću vjerojatnosti $\rho(p)$ i diskutirajte povezanost s relacijom neodređenosti.
 (c) Neka je ρ_1 gustoća vjerojatnosti nalaženja čestice s impulsom $\hbar/(4L)$. Pokažite da je

$$\frac{\rho_1}{\rho(p=0)} = \frac{4}{\pi^2}$$

8.2 Nadite valnu funkciju iz impulsnog prostora $\Phi_n(p, t)$ za n -to stacionarno stanje za problem beskonačne, pravokutne potencijalne jame širine a . Konstruirajte $|\Phi_n|^2$. Pokažite da je $|\Phi_n|^2$ konačan za $p = \pm n\pi\hbar/a$.

8.3 Izračunajte matrični element $\langle x' | H | \alpha \rangle$ gdje je $|\alpha\rangle$ po volji odabran ket, a H hamiltonijan

$$H = \frac{p^2}{2m} + V(x)$$

Kako će se rezultat promijeniti ako je $|\alpha\rangle$ svojstveni vektor od H ?

8.4 Napišite vremenski neovisnu Schrödingerovu jednadžbu u p -reprezentaciji.

8.5 Pokažite da je svojstvena funkcija operatora položaja x jednaka delta funkciji $\delta(x - x')$, a odgovarajuća svojstvena vrijednost x' .

8.6 Pokažite da vrijedi

$$\langle p' | x | \alpha \rangle = -\frac{\hbar}{i} \frac{\partial}{\partial p'} \langle p' | \alpha \rangle$$

8.1

(a) Tražimo valnu funkciju iz impulsnog prostora

$$\langle p | \psi \rangle = \psi(p) = \int_{-L}^L dx \langle p | x \rangle \underbrace{\langle x | \psi \rangle}_{\frac{1}{\sqrt{2L}}} = \frac{1}{\sqrt{2L}} \int_{-L}^L dx \langle p | x \rangle$$

$$\langle p | x \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{-ipx/\hbar}$$

$$\psi(p) = \frac{1}{\sqrt{2L}} \cdot \frac{1}{\sqrt{2\pi\hbar}} \int_{-L}^L dx e^{-ipx/\hbar} = \frac{1}{2} \cdot \frac{1}{\sqrt{\pi L \hbar}} \cdot \frac{e^{-ipx/\hbar}}{-ip/\hbar} \bigg|_{-L}^L$$

$$= -\sqrt{\frac{\hbar}{\pi L}} \cdot \frac{1}{p} \cdot \frac{1}{2i} \left(e^{-ipL/\hbar} - e^{ipL/\hbar} \right)$$

$$= \sqrt{\frac{L}{\pi \hbar}} \cdot \frac{1}{\frac{pL}{\hbar}} \cdot \sin\left(\frac{pL}{\hbar}\right)$$

Gustoća vjerovatnosti: $|\psi(p)|^2 = \rho(p)$

$$\rho(p) = \frac{L}{\pi \hbar} \cdot \frac{1}{\left(\frac{pL}{\hbar}\right)^2} \sin^2\left(\frac{pL}{\hbar}\right)$$

Da li je

$$\int_{-\infty}^{\infty} \rho(p) dp = 1 \quad ?$$

$$\frac{pL}{\hbar} = \eta \quad ; \quad dp = \frac{\hbar}{L} d\eta$$

$$\int_{-\infty}^{\infty} \frac{L}{\pi \hbar} \cdot \frac{1}{\eta^2} \sin^2 \eta \cdot \frac{\hbar}{L} d\eta = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{1}{\eta^2} \sin^2 \eta d\eta = 1 \quad \checkmark$$

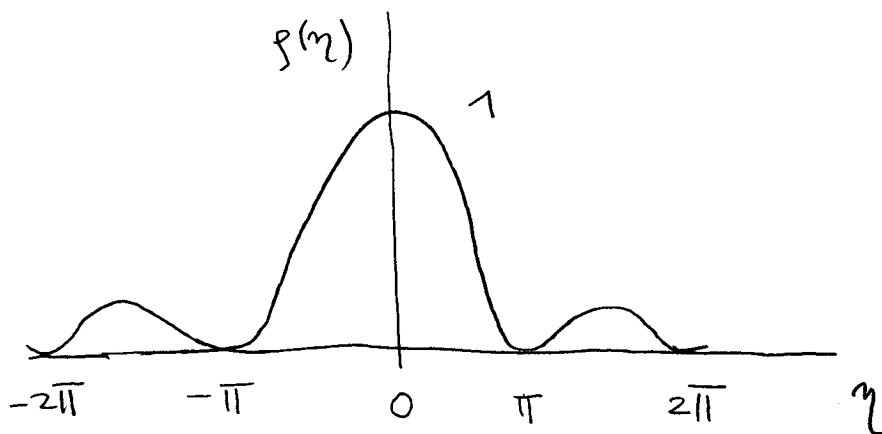
$\underbrace{\hspace{10em}}_{= \pi}$

(b) Gustočin $\rho(p)$ možemo uvesti uvedemo li

$$\eta \equiv \frac{pL}{\hbar}$$

$$\rho(\eta) = \frac{\rho(p)}{\frac{L}{\pi\hbar}} ; \rho(\eta) = \frac{1}{\eta^2} \sin^2 \eta$$

Drugim rečima, $\rho(p)$ citamo u jedinicama $\frac{L}{\pi\hbar}$ a p citamo u jedinicama $\frac{\hbar}{L}$



Neodređenost za x je

$$\Delta x \sim L$$

Neodređenost za p je

$$\Delta \eta \sim \pi \Rightarrow \Delta p \sim \pi \cdot \frac{\hbar}{L}$$

Imamo

$$\Delta x \cdot \Delta p \sim L \cdot \pi \frac{\hbar}{L} = \pi \hbar > \hbar$$

$$\Delta x \cdot \Delta p \geq \hbar$$

Relacija neodređenosti za x i p je zadovoljena!

(c)

$$\rho_1 = \rho\left(\frac{\hbar}{4L}\right) = \frac{L}{\pi\hbar} \cdot \frac{1}{\left(\frac{\hbar}{4L} \cdot \frac{L}{\hbar}\right)^2} \sin^2\left(\frac{\hbar}{4L} \cdot \frac{L}{\hbar}\right)$$

$$S_1 = \frac{L}{\pi h} \cdot \frac{1}{\left(\frac{\pi}{2}\right)^2} \underbrace{\sin^2\left(\frac{\pi}{2}\right)}_{=1} = \frac{4L}{\pi^3 h}$$

$$f(0) = \lim_{p \rightarrow 0} \frac{L}{\pi h} \cdot \frac{1}{\left(\frac{pL}{h}\right)^2} \cdot \sin^2\left(\frac{pL}{h}\right) = \frac{L}{\pi h}$$

Quanto

$$\frac{S_1}{f(0)} = \frac{4}{\pi^2}$$

Valne funkcije za beskonacnim potencijalom jamu širine L

$$\psi_n(x,t) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi}{L}x\right) e^{-iE_n t/\hbar}$$

Valne funkcije za beskonacnim potencijalom jamu iz impulsnog prostora

$$\Phi_n(p,t) = \frac{1}{\sqrt{2\pi\hbar}} \int_0^L dx \psi_n(x,t) e^{-i\frac{p \cdot x}{\hbar}}$$

$$t=0; \quad \Phi_n(p) = \frac{1}{\sqrt{2\pi\hbar}} \cdot \sqrt{\frac{2}{L}} \int_0^L dx \sin\left(\frac{n\pi}{L}x\right) e^{-i\frac{p \cdot x}{\hbar}}$$

$$\left(\frac{e^{i\frac{n\pi}{L}x} - e^{-i\frac{n\pi}{L}x}}{2i} \right)$$

$$= \sqrt{\frac{1}{\pi\hbar L}} \cdot \frac{1}{2i} \int_0^L dx \left\{ e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)x} - e^{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)x} \right\}$$

$$\int_0^L dx e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)x} = \frac{1}{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)} e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)x} \Big|_0^L$$

$$= \frac{1}{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)} \left\{ e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)L} - 1 \right\}$$

$$\int_0^L dx e^{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)x} = \frac{1}{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)} \left\{ e^{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)L} - 1 \right\}$$

$$\Phi_n(p) = \frac{1}{\sqrt{\pi\hbar L}} \cdot \frac{1}{2} \cdot \left\{ \frac{1}{\left(\frac{p}{\hbar} - \frac{n\pi}{L}\right)} \left[e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)L} - 1 \right] \right.$$

$$\left. - \frac{1}{\left(\frac{p}{\hbar} + \frac{n\pi}{L}\right)} \left[e^{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)L} - 1 \right] \right\}$$

$$\Phi_n(p) = \frac{1}{2} \cdot \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{\left(\frac{p}{\hbar}\right)^2 - \left(\frac{n\pi}{L}\right)^2} \cdot \left\{ \left(\frac{p}{\hbar} + \frac{n\pi}{L}\right) \left(e^{i\left(\frac{n\pi}{L} - \frac{p}{\hbar}\right)L} - 1\right) - \left(\frac{p}{\hbar} - \frac{n\pi}{L}\right) \left(e^{-i\left(\frac{n\pi}{L} + \frac{p}{\hbar}\right)L} - 1\right) \right\}$$

$$e^{i \frac{n\pi}{L} \cdot L} = e^{in\pi} = (-1)^n$$

$$\begin{aligned} \Phi_n(p) &= \frac{1}{2} \cdot \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{\left(\frac{p}{\hbar}\right)^2 - \left(\frac{n\pi}{L}\right)^2} \cdot \left\{ (-1)^n e^{-i \frac{pL}{\hbar}} \left(\frac{p}{\hbar} + \frac{n\pi}{L}\right) - (-1)^n \left(\frac{p}{\hbar} - \frac{n\pi}{L}\right) e^{-i \frac{pL}{\hbar}} - \frac{2n\pi}{L} \right\} \\ &= \frac{1}{2} \cdot \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{\left(\frac{p}{\hbar}\right)^2 - \left(\frac{n\pi}{L}\right)^2} \cdot \left\{ (-1)^n e^{-i \frac{pL}{\hbar}} \cdot \frac{2n\pi}{L} - \frac{2n\pi}{L} \right\} \\ &= \frac{n\pi}{L} \cdot \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{\left(\frac{p}{\hbar}\right)^2 - \left(\frac{n\pi}{L}\right)^2} \cdot \left\{ (-1)^n e^{-i \frac{pL}{\hbar}} - 1 \right\} \end{aligned}$$

$$\Phi_n(p, t) = \frac{n\pi}{L} \cdot \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{\left(\frac{p}{\hbar}\right)^2 - \left(\frac{n\pi}{L}\right)^2} \cdot \left\{ (-1)^n e^{-i \frac{pL}{\hbar}} - 1 \right\} e^{-i \frac{E_n t}{\hbar}}$$

Ako je $p = \frac{n\pi}{L} \hbar$, tada

$$\begin{aligned} \Phi_n(p) &= \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{2i} \int_0^L dx \left\{ 1 - e^{-i \frac{2n\pi}{L} x} \right\} = \\ &= \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{2i} \left\{ L - \frac{e^{-i \frac{2n\pi}{L} x}}{\left(-\frac{2n\pi}{L}\right)} \Big|_0^L \right\} \\ &= \frac{1}{\sqrt{\pi \hbar L}} \cdot \frac{1}{2i} \left\{ L + \frac{L}{2n\pi} \cdot (1 - 1) \right\} = \frac{L}{2i} \cdot \frac{1}{\sqrt{\pi \hbar L}} \\ &= \frac{1}{2\pi i} \cdot \sqrt{\frac{L}{\hbar}} \rightarrow \text{konačno pa je i } |\Phi_n(p)|^2 \text{ konačno} \end{aligned}$$

$$S_n(P) = |\Phi_n|^2$$

$$= \frac{n^2 \pi^2}{L^2} \cdot \frac{1}{\pi k L} \cdot \frac{1}{\left[\left(\frac{P}{k}\right)^2 - \left(\frac{n\pi}{L}\right)^2\right]^2} \cdot \left\{ (-1)^n e^{-i \frac{PL}{k}} - 1 \right\} \\ \cdot \left\{ (-1)^n e^{i \frac{PL}{k}} - 1 \right\}$$

$$= \frac{n^2 \pi}{k L^3} \cdot \frac{1}{\left[\left(\frac{P}{k}\right)^2 - \left(\frac{n\pi}{L}\right)^2\right]^2} \cdot \left\{ 1 - \underbrace{(-1)^n \left(e^{-i \frac{PL}{k}} + e^{i \frac{PL}{k}} \right)}_{2 \cos\left(\frac{PL}{k}\right)} + 1 \right\}$$

$$= \frac{2\pi n^2}{k L^3} \cdot \frac{1}{\left[\left(\frac{P}{k}\right)^2 - \left(\frac{n\pi}{L}\right)^2\right]^2} \cdot \left\{ 1 - (-1)^n \cdot \cos\left(\frac{PL}{k}\right) \right\}$$

8.3

$$H = \frac{p^2}{2m} + V(x)$$

$$\langle x' | p | \alpha \rangle = \frac{\hbar}{i} \frac{\partial}{\partial x'} \langle x' | \alpha \rangle$$

$$\langle x' | H | \alpha \rangle = \frac{1}{2m} \langle x' | p^2 | \alpha \rangle + \langle x' | V(x) | \alpha \rangle$$

Računamo matični element $\langle x' | p^2 | \alpha \rangle$

$$\begin{aligned} \langle x' | p^2 | \alpha \rangle &= \langle x' | p I p | \alpha \rangle = \left[I = \int dx |x\rangle \langle x| \right] \\ &= \int dx \langle x' | p | x \rangle \langle x | p | \alpha \rangle \\ &= \int dx \frac{\hbar}{i} \frac{\partial}{\partial x'} \langle x' | x \rangle \frac{\hbar}{i} \frac{\partial}{\partial x} \langle x | \alpha \rangle \end{aligned}$$

$$\begin{aligned} \langle x' | x \rangle &= \delta(x - x') \quad \text{ortogonalnost vlastitih vektora} \\ \langle x | \alpha \rangle &= \psi(x) \end{aligned}$$

$$= \int dx \left(\frac{\hbar}{i} \right)^2 \frac{\partial}{\partial x'} \delta(x - x') \frac{\partial \psi}{\partial x}$$

$$\frac{\partial}{\partial x'} \delta(x - x') = \delta'(x - x') \cdot (-1) = - \frac{\partial}{\partial x} \delta(x - x')$$

$$\text{Svojstvo datе funkcije: } \int \frac{d}{dx} \delta(x - x') f(x) dx = -f'(x')$$

$$= - \hbar^2 \frac{\partial^2 \psi}{\partial x'^2}$$

Računamo matični element $\langle x' | V(x) | \alpha \rangle$. Znamo da je

$$x |x'\rangle = x' |x'\rangle$$

Ako je $V(x)$ analitička vrijedi

$$V(x) |x'\rangle = V(x') |x'\rangle$$

$$\begin{aligned} \langle x' | V(x) | \alpha \rangle &= (\langle x' | V^\dagger(x)) \cdot | \alpha \rangle = V(x') \langle x' | \alpha \rangle \\ &= V(x') \psi(x') \end{aligned}$$

Dokazimo

$$\langle x' | H | \alpha \rangle = \frac{-\hbar^2}{2m} \frac{\partial^2 \psi(x')}{\partial x'^2} + V(x') \psi(x')$$

što smo, uobičajeno u valnoj mehanici, pisali kao

$$H \psi$$

Ako je $|\alpha\rangle$ svojstveni vektor od H

$$H |\alpha\rangle = \lambda |\alpha\rangle$$

S druge strane je

$$\begin{aligned} \langle x' | H | \alpha \rangle &= \langle x' | \cdot (H |\alpha\rangle) = \langle x' | \cdot (\lambda |\alpha\rangle) \\ &= \lambda \langle x' | \alpha \rangle \\ &= \lambda \psi(x') \end{aligned}$$

Sve skupa

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x')}{\partial x'^2} + V(x') \psi(x') = \lambda \psi(x')$$

Što je jednačina svojstvenih vrijednosti za H i $\psi(x')$ je svojstvena funkcija.

8.4

Označimo valnu funkcijo iz impulznog prostora s

$$\Phi(p) = \langle p | \alpha \rangle$$

Ako je $|\alpha\rangle$ vlatni vektor od H

$$H|\alpha\rangle = \lambda|\alpha\rangle$$

tada je $\Phi(p)$ vlatna funkcija od H u p -reprezentaciji.

Računamo matricni element $\langle p' | H | \alpha \rangle$. Na jednoj strani

$$\begin{aligned} \langle p' | H | \alpha \rangle &= \langle p' | (H | \alpha \rangle) = \lambda \langle p' | \alpha \rangle \\ &= \lambda \Phi(p') \end{aligned}$$

Na drugoj strani

$$\langle p' | \left(\frac{p^2}{2m} + V(x) \right) | \alpha \rangle = \frac{1}{2m} \langle p' | p^2 | \alpha \rangle + \langle p' | V(x) | \alpha \rangle$$

Računamo matricni element $\langle p' | p^2 | \alpha \rangle$

$$\langle p' | p^2 | \alpha \rangle = (\langle p' | (p^\dagger)^2 \rangle \cdot | \alpha \rangle) = p'^2 \langle p' | \alpha \rangle$$

jer je

$$p | p' \rangle = p' | p' \rangle \quad i \quad \langle p' | p^\dagger = p' \langle p' |$$

$$\langle p' | V(x) | \alpha \rangle = \langle p' | I \cdot V(x) \cdot I | \alpha \rangle$$

$$= \int dx' \int dp'' \langle p' | x' \rangle \langle x' | V(x) | p'' \rangle \langle p'' | \alpha \rangle$$

$$\langle p' | x' \rangle = (\langle x' | p' \rangle)^* = \frac{1}{\sqrt{2\pi\hbar}} e^{-i \frac{p' x'}{\hbar}}$$

$$\langle x' | V(x) | p'' \rangle = (\langle x' | V(x) \rangle \cdot | p'' \rangle) = V(x') \langle x' | p'' \rangle$$

$$\langle p' | V(x) | \alpha \rangle = \frac{1}{2\pi\hbar} \int dx' \int dp'' e^{-i \frac{(p' - p'') \cdot x'}{\hbar}} V(x') \Phi(p'')$$

Fourierov transformat od $V(x')$

$$V(p'-p'') = \frac{1}{\sqrt{2\pi\hbar}} \int dx' e^{-i \frac{(p'-p'') \cdot x'}{\hbar}} V(x')$$

Stoga je

$$\langle p' | V(x) | \alpha \rangle = \frac{1}{\sqrt{2\pi\hbar}} \int dp'' V(p'-p'') \phi(p'')$$

Dobijemo

$$\frac{p'^2}{2m} \phi(p') + \frac{1}{\sqrt{2\pi\hbar}} \int dp'' V(p'-p'') \phi(p'') = \lambda \phi(p')$$

što je traženi rezultat.

8.5

Operator položaja $\hat{x}$ djeluje na valnu funkciju tako da je pomnoži s x

$$\hat{x} \psi(x) = x \psi(x)$$

Ovo možemo zapisati pomoću ket-bra notacije ovako:

$$\begin{aligned} \langle x | \hat{x} | \alpha \rangle &= (\langle x | \hat{x}) \cdot | \alpha \rangle = (x \langle x |) \cdot | \alpha \rangle \\ &= x \langle x | \alpha \rangle \end{aligned}$$

gdje je $|x\rangle$ sopstveni vektor, a x svojstvena vrijednost za $\hat{x}$. Budući je $\psi(x) \equiv \langle x | \alpha \rangle$ vidimo da je

$$\hat{x} \psi(x) \leftrightarrow \langle x | \hat{x} | \alpha \rangle$$

Vektor $|\alpha\rangle$ je po volji odabran vektor stanja, pa odaberemo

$$|\alpha\rangle = |x'\rangle$$

Imamo

$$\begin{aligned} \langle x | \hat{x} | x' \rangle &= \langle x | (\hat{x} | x' \rangle) = \langle x | (x' | x' \rangle) \\ &= x' \langle x | x' \rangle = x' \delta(x - x') \end{aligned}$$

Posljednja jednakost sledi iz ortogonalnosti sopstvenih vektora $|x\rangle$ i $|x'\rangle$. Dobivamo:

$$\begin{aligned} \langle x | \hat{x} | x' \rangle &\leftrightarrow \hat{x} \delta(x - x') \\ &\parallel \\ x' \delta(x - x') \end{aligned}$$

Sve skupa,

$$\hat{x} \delta(x - x') = x' \delta(x - x')$$

pa je $\delta(x - x')$ sopstvena funkcija za $\hat{x}$ sa svojstvenom vrijednošću x' .

8.6

Koristit ćemo relacije potpunosti

$$\int dp'' |p''\rangle \langle p''| = 11$$

$$\int dx' |x'\rangle \langle x'| = 11$$

$$\begin{aligned} \langle p'|x|\alpha\rangle &= \int dp'' \langle p'|x|p''\rangle \langle p''|\alpha\rangle \\ &= \int dp'' \int dx' \langle p'|x|x'\rangle \langle x'|p''\rangle \langle p''|\alpha\rangle \end{aligned}$$

No, jednadžba svojstvenih vrijednosti za operator položaja x daje

$$x|x'\rangle = x'|x'\rangle$$

Također,

$$\langle x'|p''\rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ix'p''/\hbar}$$

Imamo

$$\begin{aligned} \langle p'|x|\alpha\rangle &= \int dx' \int dp'' \frac{x'}{2\pi\hbar} e^{-ix'p'/\hbar} \cdot e^{ix'p''/\hbar} \underbrace{\langle p''|\alpha\rangle}_{\Phi(p'')} \\ &= \frac{1}{2\pi\hbar} \int dp'' \Phi(p'') \int dx' \frac{\partial}{\partial p''} \left[e^{ix'(p''-p')/\hbar} \cdot (-i)\hbar \right] \end{aligned}$$

Druga dva člana funkcije deflirane je nula

$$\int_{-\infty}^{\infty} dx' e^{ix'(p''-p')/\hbar} = 2\pi\hbar \delta(p''-p')$$

Truando

$$\langle p' | x | \alpha \rangle = \frac{\hbar}{i} \int dp'' \Phi(p'') \frac{d}{dp''} \delta(p'' - p')$$

Calcoliamo usando la proprietà delta-funzione

$$\int f(x) \frac{d}{dx} \delta(x - x') dx = - \frac{d}{dx'} f(x')$$

Prima parte,

$$\langle p' | x | \alpha \rangle = -\frac{\hbar}{i} \frac{\partial}{\partial p'} \Phi(p')$$